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# ESTIMATION OF THE COVALENT BINDING FACTORS IN VO<sup>2+</sup> DOPED SINGLE CRYSTALS

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## **ABSTRACT**

The Steven's model is made to interpret the g-factors in various single crystals doped with paramagnetic ion  $VO^{2+}$ . This model has been used for the evaluation of covalent binding factors  $K_{\parallel}$  and  $K_{\perp}$  in different Vanadyl ions complexes with tetragonal symmetry of crystalline field. Using cubic nature for the crystalline field and covalent binding, g-factors are determined in the terms of binding parameters  $K_{\parallel}$  and  $K_{\perp}$  with the help of self developed computer program. Computations of binding factors show that one of the covalent binding parameters  $K_{\perp}$  should be less than 0.095 in order to fit the best experimental g-values.

# **KEYWORDS:**

: Electron Paramagnetic Resonance; binding parameters; spin-orbit interaction;

Fermi contact term; hyperfine interaction parameters; parametric angle.

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# INTRODUCTION

Electron Paramagnetic Resonance (EPR) studies of vanadyl ion in different single—crystals has been reported earlier [1-7] in order to study the structure as well as the relation between temperature and structure of vanadyl complexes in various host lattices. Vanadyl ion is found to be most stable cation among the few molecular paramagnetic transition metal ions, therefore a large number of paramagnetic resonance absorption studies as well as optical and magnetic studies have been done in the complexes doped with VO<sup>2+</sup> ion. Due to its higher stability, vanadyl ion is used extensively as an impurity probe to study defect properties, structural changes and host lattice effects on molecular ion using EPR studies [8]. Vanadyl complexes have been the subject of interest to a number of workers over recent years [9-10] in different crystalline field environments. The different earlier research works [1-5] show that VO<sup>2+</sup> ion always occurs co-ordinated to other groups whether in the solid state or in the solution form. The electronic state of this ion is highly dependent on the 3d<sup>1</sup> electron of vanadium. This is why the energy levels of VO<sup>2+</sup> ion can be taken as that of the V<sup>4+</sup> ion. The d<sup>1</sup> configuration of Ti<sup>3+</sup> and d<sup>9</sup> configuration of Cu<sup>2+</sup> is similar to that of the single unpaired electron vanadyl ion linked to the

 $V^{4+}$  ion. The V-O bond has larger covalent character in vanadyl ion [11]. This covalency is inversely proportional to the EPR parameters  $g_{\parallel}$  and  $g_{\perp}$ . Therefore, it is interesting to see whether it is possible to explain the experimental data obtained from the EPR spectra of  $VO^{2+}$  ion doped in different diamagnetic lattices as well as to find the bonding character of vanadium complexes on the basis of theoretical analysis of the behavior of a single  $3d^1$  electron.

It was suggested by Gerritsen and Lewis [12] that the simple expressions of g- factors in the parallel and perpendicular orientations as given by Bleaney et al. (1955) [13] could not explain the experimental g values due the presence of the crystalline field of lower symmetry and the factors not taken into analysis by Bleaney et al (1955) [13]. Steven's model [14] for covalent bonding seems to be a better tool for the explanation of experimental g-values in the presence of tetragonal crystalline field. Using experimental g-factors. In the present work it is tried to observe whether the g-values determined from the EPR experiment and also the bonding character for  $VO^{2+}$  ion in different single crystals can be explained by Steven's model (1953) [14] of covalent bonding. It is clear from the different theoretical observations that in order to have better fit the experimental g-factors one of the covalent binding factors should be taken as less than one. In table 1the covalent binding factors  $K_{\parallel}$  and  $K_{\perp}$  along with the parametric angle  $\delta$  are given. In my study for best fit to the observed g-values  $K_{\perp}$  should be less than 0.095. The following section gives the theoretical expressions of g-factors for covalent binding in the presence of tetragonal field. The discussion will be given in the subsequent section.

#### **THEORY**

In this section of the paper the theoretical expressions for g-factors are given in the presence of a tetragonal crystalline field and covalent binding for the present work. The expressions for  $g_{\parallel}$  and  $g_{\perp}$  for  $(d\epsilon)^5$  in strong cubic field having some tetragonal component were given by Steven's [14] by taking  $(d\epsilon)^5$  as a positive hole in the d $\epsilon$  shell. For the present work these expressions are slightly modified. These formulas are also used to interpret the g-factors of different single crystals.

In most of the cases the symmetry of the crystalline field is predominantly cubic, with small distortions. The splitting in the orbital states because of the cubic field is larger than that due to the terms with lower symmetry. In my case, the non-cubic component of the crystalline field consists of a tetragonal component which can be seen from observed EPR spectra. First of all, using the first order perturbation theory the matrix of cubic field splitting is diagonalized and thereafter the effect of tetragonal symmetry terms is considered. The "spin-operator technique" [15-16] is used for estimating the energy of lower states.

In the presence of a cubic field with a tetragonal distortion the equivalent Hamiltonian operator is given by:

$$\mathcal{H} = B_4(O_4^0 + 5O_4^4) + B_2^0 O_2^0 + B_4^0 O_4^0 \tag{1}$$

Here  $B_4$  is the magnitude of cubic field; second and third terms represent the tetragonal distortion of the second and fourth order in the potential, respectively. The energy levels and wave functions for the lower states are given below:

State
Doublet 
$$\begin{vmatrix} 2^a \\ 2^a \end{vmatrix} \equiv \begin{vmatrix} \widetilde{O} \\ 0 \end{vmatrix}$$

Singlet  $\pm |\mp 1\rangle \equiv |\pm \widetilde{1}\rangle$  (2a)

Energy
$$-3B_2^0 - 48B_4^0$$

$$6B_2^0 + 12B_4^0$$
 (2b)

States are given in terms of  $|L_z\rangle$  and  $|\widetilde{I}_z\rangle$  where  $\widetilde{I}$  is the fictitious orbital momentum. According to Stevens [14] the states  $|1\rangle$  and  $|-1\rangle$  do not show a purely d-nature due to  $\pi$ - bonding between d-electrons and the surrounding p-electrons of the environment. Due to the combined effect of the tetragonal field with negative  $\Delta$  and spin-orbit interaction two lower states are obtained as below:

$$\Psi_{1} = Cos \delta \left| -1 \right\rangle \left| -\frac{1}{2} \right\rangle - Sin \delta \left| 0 \right\rangle \left| \frac{1}{2} \right\rangle$$

$$\Psi_{2} = Cos \delta \left| 1 \right\rangle \left| \frac{1}{2} \right\rangle - Sin \delta \left| 0 \right\rangle \left| -\frac{1}{2} \right\rangle$$
(3)
Where,  $\tan 2\delta = \frac{\sqrt{2\lambda_{1}}}{(\lambda_{1}/2) + \Delta}$ 

and  $\lambda_1$  is the reduced value for spin-orbit interaction with  $\Delta$  as the energy difference between the states  $\Gamma_5$  and  $\Gamma_3$ . The g-factors are obtained as

$$g_{\parallel} = 2 \left| Sin^{2} \delta - (1 + \kappa) Cos^{2} \delta \right|,$$

$$g_{\perp} = 2 \left| \sqrt{2 \kappa} Sin \delta Cos \delta + Sin^{2} \delta \right|$$
(5)

Where K is defined as  $K = \langle a | L_Z | a \rangle$  and it is assumed that covalent binding is isotropic, that is, same in all directions. If covalent binding is anisotropic, then expressions for g-factors will be as:

$$g_{\parallel} = 2 \left| Sin^{2} \delta - (1 + \kappa_{\parallel}) Cos^{2} \delta \right|,$$

$$g_{\perp} = 2 \left| \sqrt{2} \kappa_{\perp} Sin \delta Cos \delta + Sin^{2} \delta \right|$$
Where  $\kappa_{\parallel} = \left\langle 1 \middle| L_{z} \middle| 1 \right\rangle$  and  $\kappa_{\perp} = \sqrt{2} \left\langle 1 \middle| L_{z} \middle| 0 \right\rangle$ 

# RESULTS AND DISCUSSION

In general, the EPR spectra of VO<sup>2+</sup> in single crystal are described by a strong octahedral field having tetragonal distortion due to covalent bonding between the central metal ion and surrounding ligand. Using equation (6) the covalent binding factors  $K_{||}$  and  $K_{\perp}$  were estimated. The values of these factors for different single crystals [17-34] with parametric angle  $\delta$  satisfying experimental g-values are given in Table 1. In presence of tetragonal field, the covalent binding factors are found to be less than one, because no shifting of doublet in  $t_{2g}$  and higher doublet  $e_g$ . The multiple values of [35] of  $K_{||}$ ,  $K_{\perp}$  and  $\delta$  have been estimated as exact value of  $\lambda_1$  is not known. In case of exact value parametric angle  $\delta$  can be fixed so that corresponding

parameters with respect to  $\delta$  is obtained. The theoretical calculations show that experimental values of g-factors can be obtained if  $\mathbf{K}_{\perp}$  is less than 0.095. The larger migration of an electron to its neighboring atom may be explained by less value of covalent binding factor [36]. However, covalency is inversely proportional to  $\mathbf{g}_{\parallel}$  and  $\mathbf{g}_{\perp}$ . Also,  $\mathbf{g}_{\parallel}$  varies inversely to covalent binding factor  $\mathbf{K}_{||}$ , that is, if  $\mathbf{K}_{||}$  is decreased  $\mathbf{g}_{\parallel}$  will be increase and so covalency will decrease indicating more ionic character of the complex [37]. The more ionic the crystal is larger will be splitting [38-39]. However direct proportionality between  $\mathbf{g}_{\perp}$  and  $\mathbf{K}_{\perp}$  will give a different nature to the complex as for as binding is considered. The rate of decrease of  $\mathbf{K}_{||}$  is more than  $\mathbf{K}_{\perp}$  and therefore taking more ionic nature for the complex will be appropriate. Therefore, the migration of electron to its neighboring atoms will take place. These assumptions indicate that crystal field theory for ionic complex must be modified by considering a definite amount of charge transfer.

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Table 1 Values of the covalent bonding factors  $\,K_{\,\parallel}\,\text{and}\,\,K_{\,\perp}\,$  with parametric angle  $\delta$  in different crystals

S. No.	Crystal Lattices	$K_{\perp}$	К	δ	Reference
1.	Ammonium Alum	0.030	1.629	95°45′	[17]
1.	{NH <sub>4</sub> Al (SO <sub>4</sub> ) <sub>2</sub> .24H <sub>2</sub> O}	0.035	0.988	96°	[1/]
	{1V114A1 (5O4)2.24112O }	0.040	0.662	96°15′	
		0.047	0.455	96°30′	
		0.053	0.239	96°45′	
		0.062	0.239	90 43 97°	
2.	Rb Alum	0.010	1.145	96°14′	[18]
	{Rb Al (SO <sub>4</sub> ) <sub>2</sub> . 12H <sub>2</sub> O}	0.018	0.889	96°49′	
	[ - ( 1/2)	0.028	0.689	97°30′	
		0.034	0.339	97°52′	
		0.037	0.074	98°59′	
					54.03
3.	L-asparagine monohydrate	0.022	1.485	96°14′	[19]
		0.010	0.933	96°49′	
		0.007	0.833	97°30′	
		0.006	0.736	97°52′	
		0.041	1.324	95°45′	
4.	Sodium Citrate	0.033	0.966	93 43 96°	[20]
••		0.020	0.604		[=0]
		0.010	0.324	96°15′	
		0.007	0.324	96°30′	
		0.004	0.243	96°45′ 97°	
		0.031	1.621		
5.	Tutton Salts	0.040	1.631	96°28′	[21]
٥.	Tutton Saits	0.040	0.985	96°44′	[21]
		0.047	0.663	97°14′	
			0.495	97°32′	
		0.057	0.215	97°46′	
		0.062	0.061	98°2′	
		0.061	1 140	95°24.5′	
6.	Tri Sodium Citrate Pentahydrate	0.063	1.148		[22]
	${Na_3C_6H_5O_7.5H_2O}$	0.068	0.951	95°48′	
		0.072	0.612	95°48′	
		0.061	0.319 0.074	96°14′ 96°32′	
		0.031	1.174		
		0.037	0.898	96°49′	[23]
7.	Calcium Fumarate Trihydrate	0.037	0.898	97°8′	
		0.047		97°20′	
		0.047	0.357	97°40.5′	
		0.061	0.162	97°54.5′	
		0.001	0.041	98°28′	
_	CAMTH	0.030	1.173	96°29′	[0.4]
8.	CoMTH	0.036	0.897	96°47.5′	[24]
	Cobalt Maleate Tetrahydrate			70 41.3	

				<del></del>	
		0.040	0.729	96°59.5′	
		0.050	0.356	97°32′	
		0.056	0.161	97°46′	
		0.060	0.042	98°2′	
9.	Potassium Oxlate Monohydrate (POM)	0.004	0.585	96°35′	[25]
· ·	(COOK) <sub>2</sub> .H <sub>2</sub> O	0.006	0.522	96°42′	[20]
	(6001)2.1120	0.007	0.461	96°48′	
		0.012	0.397	96°55.5′	
		0.012	0.397		
		0.020	0.110	97°16′	
10	Y CONTROL OF CONTROL	0.001	0.918	96°11.5′	F2 61
10.	Lactate Trihydrate (ZLT)	0.002	0.804	96°15.5′	[26]
	[Zn (CH <sub>3</sub> CHOHCOO) <sub>2</sub> .3H <sub>2</sub> O]	0.004	0.660	96°29.5′	
		0.008	0.595	96°35.5′	
		0.010	0.277	96°46′	
		0.020	0.002	97°2′	
		0.002	0.477	95°23′	
11.	KHC <sub>2</sub> O <sub>4</sub>	0.004	0.413	95°27.5′	[27]
		0.006	0.353	95°32.5′	
		0.008	0.333	95°37′	
		0.010	0.234	95°43.5′	
		0.001	0.688	96°29′	
		0.004			[28]
12.	TiO <sub>2</sub>	0.004	0.648	96°47.5′	[26]
12.	1102	0.008	0.584	96°59.5′	
		0.010	0.445	97°32′	
			0.384	97°46′	
		0.025	0.087	98°	
		0.010	1.047	0.5024	
			1.047	96°2′	[29]
13.	GeO <sub>2</sub> (amorphous)	0.015	0.798	96°16′	[29]
13.	GeO <sub>2</sub> (amorphous)	0.020	0.642	96°29.5′	
		0.025	0.424	96°44.5′	
		0.031	0.292	96°56′	
		0.036	0.108	97°13′	
		0.010	1.187	96°1′	
14.	FeRb <sub>2</sub> (SO4) <sub>2</sub> .6H <sub>2</sub> O (Tutton Salt)	0.016	0.766		[30]
		0.022	0.537	96°17′	
		0.028	0.398	96°59.5′	
		0.036	0.398	97°32′ 97°46′	
15.	NaKC <sub>4</sub> H <sub>4</sub> O <sub>6</sub> .6H <sub>2</sub> O	0.001	0.914	96°36′	
		0.003	0.804	96°41′	[30]
		0.005	0.660	96°45′	
		0.008	0.660	96°51.5′	
		0.011	0.392	97°16′	
		0.020	0.274	97°20.5′	
		0.060	1.447	94°52′	
16.	MgSO <sub>4</sub> .7H <sub>2</sub> O	0.063	1.447	95°0.5′	[30]
		0.066	0.948	95°19.5′	
	1		U.948		
		0.071	0.604	95°38′	

		0.075	0.317	95°57.5′	
		0.079	0.076	96°23.5′	
17.	MASH	0.060 0.066 0.070 0.075 0.080 0.085	1.027 0.889 0.687 0.492 0.227 0.074	96°23′ 96°47.5′ 97°3′ 97°28′ 97°46′ 98°2.5′	[31]
18.	Cs <sub>2</sub> Mg (SO4) <sub>2</sub> .6H <sub>2</sub> O	0.030 0.040 0.045 0.051 0.056 0.060	1.628 0.980 0.660 0.468 0.215 0.062	96°15′ 96°47.5′ 97°8.5′ 97°22′ 97°42′ 97°57′	[32]
19.	Mg (H <sub>2</sub> O)6H <sub>2</sub>	0.066 0.069 0.071 0.075 0.080 0.082	1.066 0.928 0.807 0.483 0.295 0.207	96°21.5′ 96°30′ 96°39.5′ 97°3′ 97°20′ 97°29′	[32]
20.	CoASH	0.030 0.034 0.039 0.045 0.052 0.060	1.273 0.947 0.788 0.456 0.199 0.052	98°46′ 99°7′ 99°21.5′ 99°54.5′ 100°16′ 100°32′	[33]
21.	CASH	0.030 0.040 0.050 0.070 0.080 0.090	1.153 0.818 0.616 0.383 0.284 0.194	96°21′ 96°29.5′ 96°38.5′ 97°3′ 97°20′ 98°28′	[34]

### **REFERENCES:**

- [1] G. M. Zverev and A. M. Prokhorov, Sov. Phys.-JETP 12 (1960) 160.
- [2] M. B. Plama, Vittorelli, M. V. Plama, D. Palumbo and F. Sgarlata, *Nuvo Cimento* 3 (1956) 718.
- [3] R. H. Borcherts and C. Kikuchi, J. Chem. Phys. 40 (1964) 2270.
- [4] P. T. Manoharan and M. G. Rogers, J. Chem. Phys. 49 (1968) 3912.
- [5] R. P. Dodge, D. H. Templeton and A. Zalkin, J. Chem. Phys. 35 (1961) 55.
- [6] K. de Armond, B. B. Garrett and H. S. Gutowsky, J. Chem. Phys. 42 (1965) 1019.
- [7] A. Moonigan and J. A. Mackinnon, Can. J. Phys. 45 (1967) 2769.
- [8] J. E. Wertz, J. R. Bolton, Electron Spin Resonance: Elementary Theory and practical applications, Mc Graw-Hill, New York, 1972.
- [9] T. Kawano and K. Furukrua, J. Phys. Soc. Jpn. 47 (1979) 1119.
- [10] S. Radhakrishna and M. Salagram, Solid State Commun. 47 (1983) 77.
- [11] R. P. Kohin (1979) *Magnetic Resonance Review* (ed.) C. P. Poole (New York: Gordan and Breach) Vol. 5 p.75
- [12] H. T. Gerritsen and H. R. Lewis, *Phys. Rev.* 119 (1960) 1010.
- [13] B. Bleaney, G. S. Bogle, A. H. Cooke, R. T. Duffis, M. C. M. O'Brien and K. W. H. Steven's *Proc. Phys.Soc.* A68 (1955) 57.
- [14] K. W. H. Steven's Proc. R. Soc. A219 (1953) 542.
- [15] B. Bleaney, H. E. D. Scovil and R. S. Trenam, Proc. R. Soc. A223 (1954) 15.
- [16] J. M. Baker, F. I. B. Williams, Proc. Phys. Soc. 78 (1961) 1340.
- [17] K. V. S. Rao, M. D. Sastry and P. Venkateswarlu, J. Chem. Phys. 49 (1968) 4984.
- [18] A. Manoogian and J. A. MacKinnon, Can. J. Phys. 45 (1967) 2769.
- [19] Ram Kripal and Pragya Singh, Solid State Commun. 142 (2007) 412.
- [20] Ram Kripal and Indraject Mishra, Spectrochim. Acta A 72 (2009) 538.
- [21] R. H. Borcherts and C. Kikuchi, J. Chem. Phys. 40 (1964) 2270.
- [22] Indrajeet Mishra, Int. J. Rec. Sci. Res. 8(10) (2017) 20683.
- [23] Ram Kripal, <u>Indrajeet Mishra</u>, S. K. Gupta & Manju Arora, Chem. Phys. Lett. 484 (2010) 200.
- [24] S. Radhakrishna, B. P. R. Chowdary and A.K. Viswanath J. Phys. Soc. Jpn. 41 (1976) 530.
- [25] Ram Kripal Manju Maurya and Hargovind, Physica B 392 (2007) 281.
- [26] Ram Kripal and Pragya Singh, J. Magn. Mag. Mat. 307 (2006) 308.
- [27] M. Salagram, PhD. Thesis I. I. T. Madras (1980).
- [28] H. J. Gerritsen and H. R. Lewis, Phys. Rev. 119 (1960) 1010.
- [29] I. Siegel, Phys. Rev. 134 (1964) 193A.
- [30] V. K. Jain and P. Venkateshwarlu, J. Chem. Phys. 73(1) (1979) 30.
- [31] S.K. Mishra and J. Sun, Physica B 162 (1990) 331.
- [32] M. Narayana, J. Chem. Phys 72 (1980) 4255.
- [33] S. K. Mishra, J. Sun and X. Li *Physica B* 168 (1991) 70.

- [34] S. K. Mishra and J. Sun. Physica B 162 (1990) 331.
- [35] D. K. Rei, Sov. Phys. Solid State (Engl. Transl.) 3 (1962) 1838.
- [36] C. J. Ballhausen and H. B. Gray, J. Inorg. Chem. 1 (1962) 111.
- [37] D. Kivelson and S. K. Lee, J. Chem. Phys. 41 (1964) 1869.
- [38] J. S. W. Van, Disc. Faraday Soc. 19 (1955) 19.
- [39] A. Abragam Phys. Rev. 79 (1950) 534.